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STRENGTH OF GLASS -- A FRACTURE MECHANICS APPROACH

S. M. Wiederhorn

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16. ABSTRACT (A 200 word or less factual summary of most significant information. If document includes a significant bibliography or literature survey, mention it here.)  After a brief review of those factors that determine the strength of glass (brittleness, surface flaws, susceptibility to stress corrosion cracking), a discussion will be given of how fracture mechanics techniques can be used to understand the physics and chemistry of glass strength. In this paper we assume that the strength of glass is limited by the growth of cracks that are always present in normal glass surfaces. Fracture mechanics techniques can be used to characterize the crack growth and to relate the growth to experimental parameters such as temperature, environment and glass composition. Crack growth lata obtained in this manner can be used to develop a deeper understanding of fracture mechanisms, and to develop charts that can be used for the design of glass structural components. Examples of both applications are given in the paper.				
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# STRENGTH OF GLASS--A FRACTURE MECHANICS APPROACH

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#### INTRODUCTION

In this paper fracture mechanics is discussed: as a method of understanding the strength of glass, and as a method of providing data that can be used to assure the structural reliability of glass. A review is first presented of the main factors that control the strength of glass (temperature, environment, surface condition, etc.). Then fracture mechanics techniques are described, and the available fracture mechanics data are reviewed. Finally, the use of fracture mechanics for failure prediction is discussed.

# REVIEW OF THE LITERATURE ON THE STRENGTH OF GLASS

The two main characteristics that determine the strength of glass are brittle behavior and static fatigue. Glass is one of the most frittle materials known because at low temperatures deformation is almost completely elastic. This elasticity has been demonstrated on silica rods and fibers by Hillig (Ref. 1) and by Mallinder and Proctor (Ref. 2) to loads as high 1.1 GN/m<sup>2</sup> (1.6 x 10<sup>5</sup> psi), and on soda lime silicate glass fibers by Mallinder and Proctor (Ref. 2) to 0.4 GN/m<sup>2</sup> (0.6 x 10<sup>6</sup> psi). The stresses were so high in these experiments that non-linear elastic behavior was observed; however, the deformation completely disappeared when the load was released. Interestingly, the Youngs modulus for the silica glass increased during these experiments, while that for soda lime silicate glass decreased.

Permanent deformation can occur in glass under certain circumstances. Permanent densification has been observed by Bridgeman (Ref. 3) and by Cohen and Roy (Ref. 4) on glasses that have been subjected to pressures of 7 GN/m² (10<sup>6</sup> psi). During hardness indentation tests, plastic flow and densification occurs because of the high compressive and shear stresses (~ 10 GN/m²) near the indenter (Ref. 5,6). Permanent deformation has also been observed in compressive tests conducted by Ernsberger (Ref. 7) on glass rods that contained oblate bubbles. In these experiments densification was observed in Pyrex and silica glass at stresses of approximately 7 GN/m². In all of the above cases, high compressive loads were necessary for the occurrence of plastic deformation. To date, there have been no unequivocal reports of plastic deformation of glass that has been subjected to tensile loads.

The low strength of glass is a direct result of this brittle behavior. When glass is polished and ground, small cracks are left in the glass surface. These cracks act as stress concentrators so that the crack tip stress is considerably higher than the average applied stress. Since plastic deformation does not occur to relieve these crack tip stresses, the surface cracks propagate when a critical stress is reached at the crack tip. The critical stress required for crack motion is determined in part by the glass composition, and in part by the environment at the crack tip. The environmental contribution to crack motion results in a time dependence of strength known as static fatigue.

Static fatigue of glass was observed first by Grenet (Ref. 9) who noted a time dependence of strength for both static and dynamic loading conditions. For a constant loading rate, the strength depended on the rate of loading, increasing as the rate increased; whereas for constant load, failure occurred after a period of time at load. The phenomenon of static fatigue at constant load is illustrated in figure 1.

The failure time depends on the load applied to the glass; large loads can only be supported for short periods of time. If the load is small enough, a fatigue limit occurs; fracture does not occur regardless of the time for which the load is applied.

Static fatigue can be explained by the growth of cracks caused by a stress enhanced chemical reaction between water and glass. Static fatigue occurs only in the presence of water which reacts chemically with the strained bonds at the crack tip causing bond rupture. The rate of crack growth is determined by the rate of the chemical reaction, and the time to failure is determined by the time

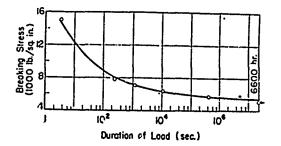


Fig. 1 Static fatigue of sodalime silicate glass (Ref. 10)

required for the crack to grow from a subcritical to a critical size, at which point, failure is instantaneous.

Before the advent of fracture mechanics concepts, strength measurements were used to study static fatigue. The strength was measured in three or four point loading as a function of loading rate, or the time to failure was measured as a function of applied load. Using these techniques the following observations were made:

- o Static fatigue occurred only in the presence of water (Ref. 11-14).
- o Static fatigue could be detected for load times as short as  $10^{-2}$  second (Ref. 14).
- o Static fatigue did not occur in a vacuum or at very low temperatures (-196°C) (Ref. 12, 13, 15).
- o Static fatigue was an activated process, fracture becoming easier at higher temperatures (Ref. 16, 17).
- o Static fatigue was the result of a single process that did not depend on the surface flaw size (Ref. 18).
- o The static fatigue limit occurs at approximately 20-30 percent of the environment-free strength (Ref. 10, 19).

While strength studies provided considerable information on the strength of glass and the occurrence of static fatigue, fracture mechanics techniques added to this body of knowledge, providing a deeper insight into the fracture process.

# FRACTURE MECHANICS TECHNIQUES

Fracture mechanics techniques are all based on Griffith's original studies of the fracture of brittle materials (Ref. 21). One conclusion of Griffith's work is that the strength of brittle materials depends on the presence of cracks in these materials. The scrength, S, was related to the crack length by the following formula (derived for an elliptical through crack of length 2a):

$$S = (2E\gamma/\pi a)^{\frac{1}{2}} \tag{1}$$

where  $\gamma$  is the fracture surface energy of the material and E is Young's modulus. Griffith showed that the fracture behavior of glass containing macroscopic cracks obeyed the above equation.

The theory developed by Griffith was generalized by Irwin (Ref. 21) who established fracture mechanics as a practical science which could be used to assure the reliability of structural materials. In Irwin's analysis the condition for failure could be expressed in the following general form for plane stress deformation:

$$X_T^2 = 2E\gamma \tag{2}$$

Where  $K_{\underline{I}}$  is known as the stress intensity factor for opening mode loading.

The stress intensity factor,  $K_{\underline{I}}$ , occupies a central position in the science of brittle fracture because it is proportional to the

stress near the crack tip. In fact at any point near the crack tip the stress,  $\sigma_{ij}$ , is related to the stress intensity factor by the following equation (Ref. 22):

$$\sigma_{i,j} = (K_T/\sqrt{r}) f_{i,j} (\theta)$$
 (3)

where r is the distance from the crack tip and  $f_{ij}$  (0) is a function of the angle 0 from the fracture plane. Eqn. (3) applies to the region near the crack tip for any crack shape, provided the crack is subjected to opening mode loading (all loads are perpendicular to the crack plane). For any particular crack  $K_{I}$  can be related to the applied load by a stress analysis of the crack geometry. Therefore, the applied load can be related to the stresses near the crack tip. For this reason,  $K_{I}$  is the main mechanical parameter that controls the fracture of brittle materials.

The relation between  $K_{\rm I}$  and the applied load has been determined for many different crack configurations (Ref. 23, 24). For a Griffith crack  $K_{\rm I}$  =  $S\sqrt{\pi a}$ , which is obtained by comparing Eqn. (1) and (2). While the Griffith configuration could be used to obtain fracture information, other crack configurations are easier to use. Two crack configurations that have been used most frequently to obtain fracture mechanics data on glass are the double cantilever beam configuration (Ref. 24, 25) and the double torsion configuration (Ref. 26). Either of these techniques can be used conveniently to obtain fracture mechanics data on glass (Ref. 27, 28).

## FRACTURE MECHANICS DATA ON GLASS

#### Crack Growth Data

The earliest fracture mechanics studies, conducted in nitrogen gas containing various amounts of water vapor (Ref. 29), showed that glass fracture was no simple process. Fracture appeared to occur by several mechanisms as indicated by the trimodal curves shown in Fig. 2. Three characteristic regions of crack growth were observed. At low values of K<sub>I</sub>, region I, the crack velocity, v, depended exponentially on K<sub>I</sub>, and also on the parital pressure of water in the environment.

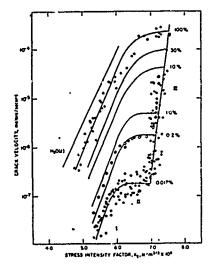


Fig. 2 Effect of water vapor (% R.H.) on crack propagation (Ref. 29)

At higher values of  $K_{\rm I}$ , region II, the crack velocity depended on the partial pressure of water in the environment, but was nearly independent of  $K_{\rm I}$ . Finally at the highest values of  $K_{\rm I}$ , region III, crack growth again depended exponentially on  $K_{\rm I}$ , but now did not depend on water in the environment. In region I and II crack growth was attributed to a chemical reaction between water in the environment and the stressed glass at the crack tip. In region I crack motion was reaction rate limited while in region II crack motion was transport rate limited depending on concentration differences between the crack tip and the bulk environment. In region III crack motion was independent of environment and fracture was due to some process that depends only on the glass structure. More recent studies of fracture in vacuum (see below) indicate that region III does not occur for all glasses.

Crack propagation in region I has been shown to depend on temperature, glass composition and environment. For constant K<sub>I</sub>, the crack velocity in water increases as the temperature is increased suggesting that crack growth is an activated process (Ref. 30). A least squares fit of crack growth data for a number of glasses to an Arhenius type equation yields a zero-stress activation energy that ranges from 20 to 30 Kcal/ mol. This range of activation energies is consistent with that expected from a chemical reaction, and is, therefore, consistent with the mechanism proposed by Charles and Hillig for the static fatigue of glass (Ref. 31, 32). In other aqueous environments (acids, bases and various salt solutions) crack propagation has been shown to depend on the pH of the test solution (Ref. 33), Fig. 3.

For example, in silica glass crack propagation curves obtained in high pH solutions have about twice the slope of those obtained in low pH solutions.

This change of slope with pH can be used to explain similar slope changes observed for other glasses that were tested in water (Ref. 30). In general, high alkali containing glasses have slopes that are more shallow than the low alkali containing glasses. This difference in slope has been attributed to a difference in pH at the crack tip, the pH being determined by a

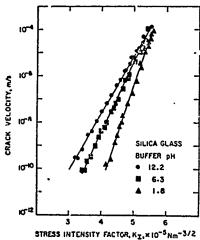


Fig. 3 Effect of pH on crack propagation (Ref. 33)

chemical reaction between the glass and the crack tip solution. The high alkali containing glasses react to form a basic crack tip environment while glasses containing little or no basic constituents react to form an acidic environment.

Region I crack growth has also been investigated in normal alcohols by Freiman (Ref. 34) and by Evans and Wiederhorn (Ref. 35). Trimodal curves were obtained (Fig. 4) indicating that crack growth behavior in these environments was similar to that in nitrogen gas containing water vapor. At a given value of K<sub>I</sub>, the crack velocity was proportional to the concentration of the water in the alcohol, suggesting that crack propagation is mainly

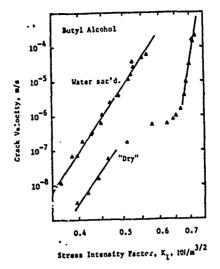


Fig. 4 Crack propagation in Butyl alcohol (Ref. 35)

due to the water in the alcohol and not to the alcohol itself. However, in region III the alcohol chain length does have a small effect in the position of the crack velocity curve.

Information on the mechanism of fracture in region III has been chained recently by Wiederhorn et al (Ref. 36) from crack propagation studies in vacuum. In these studies, it was observed that the crack growth

depended on the composition of the glass. For some glass compositions crack growth in region III depended on temperature in an Arhenian manner (Fig. 5). For other glass compositions slow crack growth did not occur; instead glass fractured abruptly at a critical value of K<sub>I</sub>. These differences in behavior apparently depend on the elastic properties of the glass. Crack growth in vacuum occurred for glasses exhibiting normal elastic behavior (the bulk modulus decreasing with increasing temperatures but increasing with increasing pressure). By contrast crack growth did not occur

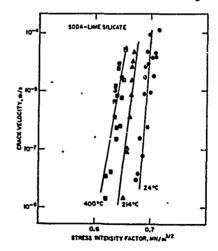


Fig. 5 Crack propagation in vacuum (Ref. 36)

for glasses that exhibited anomolous elastic behavior (the bulk modulus increasing with increasing temperatures, but decreasing with increasing pressure). Wiederhorn et al (Ref. 36) suggested that this difference in behavior could be related to the glass structure near the crack tip. Thomson's theory of "lattice trapping" which relates structure to fracture behavior may explain the fracture behavior of glass in vacuum (Ref. 37, 38). The reader is referred to the original article by Wiederhorn et al (Ref. 36) for a futher discussion of this point.

#### CRITICAL STRESS INTENSITY FACTOR DATA

In addition to crack propagation studies, fracture mechanics techniques can be used to measure the critical stress intensity factor for abrupt fracture,  $K_{\text{IC}}$ . Because  $K_{\text{T}}$  is related to the stresses and strains near the crack tip,  $K_{\text{TC}}$  gives a measurement of the maximum stresses or strains that a material can withstand prior to failure.  $K_{\text{TC}}$  is commonly defined as the value of  $K_{\mathsf{T}}$  that is required for crack growth in a inert environment.  $K_{\underline{IC}}$  is a well defined parameter for materials that fail abruptly because of rapid crack acceleration at a well defined value of  $K_{\tau}$ . However, for materials that exhibit slow crack growth in an inert environment,  $K_{TC}$  is not so easily defined because the lower limit of  $K_{T}$  for the initiation of crack motion is not easily measure. For these materials we have defined  $K_{TC}$  as the value of  $K_{T}$  (measured experimentally) required for a crack to move at a velocity  $\sim 10^{-1}$  m/s. Using this definition, values of  $K_{{\ensuremath{\mathsf{TC}}}}$  have been reported for a number of different glass compositions (Ref: 39, 40). The values of  $K_{\mbox{\scriptsize IC}}$  for glass depend on composition and range from 0.6 to 0.9  $MN/m^{3/2}$ .

The values of  $K_{\rm TC}$  for glass can be used to demonstrate the extreme brittleness of glass. Dugdale (Ref. 41) and McClintock and Irwin (Ref. 42) have shown that a good estimate of the plastic zone size, R, near a crack tip can be obtained from the following equation:

$$R = (\pi/8) (K_{IC}/\sigma_y)^2$$
 (4)

where  $\sigma_y$  is the yield stress of the glass. Applying this equation to silica glass ( $K_{IC}=0.79~\text{MN/m}^{3/2}$ ) and soda lime silicate glass ( $K_{IC}=0.75~\text{MN/m}^{3/2}$ ) for which  $\sigma_y$  has been estimated from hardness indentation studies (Ref. 43) to be 19.5 and 10 GN/m² respectively, the sizes of the plastic zones at the crack tip calculated from eqn. 4 are 6.4 x 10<sup>-10</sup>m for silica glass and 26 x 10<sup>-10</sup>m for soda lime silicate glass. The fact that the plastic zone sizes for metals and plastics (1.4 x 10<sup>-4</sup> - 14 x 10<sup>-4</sup>m for PMMA and 7.1 x 10<sup>-4</sup>m

for 4340 steel) (Ref. 44) are so much larger emphasizes the extremely brittle nature of glass. Plastic flow at crack tips in glass does not significantly reduce the stress concentrating effect of surface microcracks.

## Comparison of Crack Propagation Data with Strength Data

Crack growth data are of value for engineering purposes only if the data are consistent with strength data obtained at constant load or at constant loading rate. The easiest experimental comparise is between the crack propagation data and the loading rate data.

For most glasses the crack velocity can be expressed as a power function of  $K_T$ ,  $v = AK_T^{II}$ , where A and n are experimentally determined crack propagation parameters. Using this expression, Evans showed that the loading rate,  $\delta$ , was related to the strength,  $\sigma$ , by the following equation (Ref. 45):

$$\sigma^{n+1} = 2(n+1) \delta c_{IC}^{n-2} / Ay^2 K_{IC}^{n-2} / 2$$
 (5)

where  $\sigma_{TC}$  is the median strength of the glass measured in an inert environment (vacuum or super dry nitrogen), and Y is the factor that depends on the crack geometry. Using this equation, a comparison between crack propagation data and strain rate data was made for two glasses: a soda lime silicate glass (Ref. 47) and an ultralow expansion silica glass containing 7.5 percent TiO, (Ref. 48). Results for the silica glass tested in water (Fig. 6) show excellent agreement between strength measurement and crack propagation data suggesting that failure was primarily by crack propagation. A similar comparison for soda lime silicate glass tested in nitrogen. one percent R.H., also gives excellent agreement between the crack propagation data and

Other less complete comparisons between strength and crack propagation data have been made. In a recent studies on a variety of glasses, Ritter (Ref. 48) reports good agreement for n obtained

the strength data.

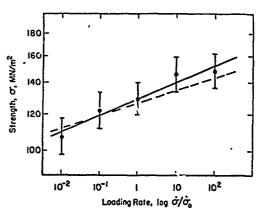


Fig. 6 A comparison of strength and crack propagation data (Ref. 48)

from loading rate and crack propagation studies. Slopes determined from constant load studies on silica and anda lime silicate glass in water are also in good agreement with slopes determined from crack propagation studies (Ref. 30). \* However, constant load studies on abraded silica glass and on Fyrex glass (abraded and chemically polished) do not agree with crack propagation studies (Ref. 49, 50). difference in behavior has been attributed to differences in crack tip chemistry for the two types of studies (Ref. 33). Thus, if the pH at the crack tip differed for the two types of studies then differences in slope would be expected. Because the failure time can be estimated from crack growth information, it is of practical importance

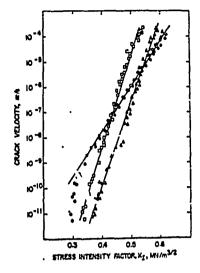


Fig. 7 Crack propagation air (50% R.H.): ∆ silica glass; □ pytex; o sodalime (Ref. 51)

# The Static Fatigue Limit in Glass

to resolve these differences in behavior.

Fracture mechanics techniques can be used to investigate the static fatigue limit in glass. This limit, which for a given component is the load below which fracture will not occur regardless of how long the component is subjected to load, has been estimated by various authors as 20 to 30 percent of the short term strength (presumably  $\sigma_{TC}$ ) (Ref. 10, 19). Fracture mechanics data can also be used to estimate the value of the static fatigue limit in glass. The fatigue limit is determined by the value of  $K_T$  at which crack motion stops. Recent studies by Wiederhorn and Johnson (Ref. 51) at crack velocities as low as  $10^{-11}$  m/s gave no indication of a static fatigue limit for Pyrex cr silica glass tested in air (50 percent R.H.). The crack velocity data for these glasses approximate a straight line on a logarithmic plot over the entire range of parameters (Fig. 7). The lowest

<sup>\*</sup>In these studies the crack velocity was expressed as an exponential function of  $K_{I}$ ,  $v = v_{O}$  exp  $\beta K_{I}$ , and the strength as  $\sigma/\sigma_{IC} = 0.5 - (1/\beta K_{IC})$ In  $(t/t_{0.5})$ , where  $t_{0.5}$  is the time to failure at a load  $\sigma_{IC}/2$ . Termed the universal fatigue curve, this type of strength plot was suggested first by Mould (Ref. 18). For glass,  $K_T$  - v data are represented equally well by an exponential or power function of  $K_T$ .

values of  $K_{\rm I}$  measured for these glasses suggest that the fatigue limit is less than  $^{\sim}44$  percent of the short term strength for both Pyrex and silica glass. By contrast, the curve for soda lime silicate glass does exhibit severe bending at low velocities suggesting a fatigue limit  $K_{\rm I} \sim 0.3$  MN/m $^{3/2}$ , which is  $^{\sim}40$  percent of the short term strength. (Defined here as  $K_{\rm I}/K_{\rm IC}$ ). This value is slightly higher than those reported in the literature. If Pyrex is tested in an acid environment (pH < 1.7) one observes a fatigue limit of  $K_{\rm I} \sim 0.41$  MN/m $^{3/2}$  which is  $^{\sim}53$  percent of the short term strength (Ref. 34). This limit results from an aggressive attack of the acid on the glass which causes crack tip rounding.

# FRACTURE MECHANICS FOR FAILURE PREDICTION

The excellent agreement between crack propagation data and strength data (Ref. 46, 47) supports the argument that failure in these glasses is by crack growth, and suggests that crack growth data can be used for failure prediction purposes. This use of crack growth data depends on a complete characterization of crack growth for a given application. Once the crack growth has been characterized then the time-to-failure can be calculated provided the size of the critical flaw can be determined. The key to this application of fracture mechanics is in fact the determination of the initial flaw size. Recent work by Evans and Wiederhorn (Ref. 52, 53) suggest that for practical purposes the initial flaw size can be estimated by overload proof testing.

Proof testing is used in structural design to break weak components before they can be placed into service. By so doing we assure the structural integrity of components that are placed in service. In a proof test, a proof test load,  $\sigma_p$ , which is larger than the service load,  $\sigma_a$ , is applied to the components in an inert environment. This procedure guarantees that components passing the proof test have flaws smaller than the critical size that would have resulted in failure during service. In fracture mechanics terms, the stress intensity factor at the crack tip of the most serious flaw,  $K_p$ , has to be less than  $K_{IC}$  if the component is not to break during the proof test. Thus, for a component to pass the proof test,

$$K_{IC} > K_{D} = \sigma_{D} Y \sqrt{a}_{i}$$
 (6)

where Y is a geometric factor and a is the crack length of the most serious flaw in the component.

When a component has passed the proof test and is placed into service the stress intensity factor at the most serious flaw is given by:

$$K_{i} = \sigma_{a} Y \sqrt{a}_{i} \tag{7}$$

where  $\sigma_a$  is the service load.

By dividing Eqn. (6) into Eqn. (7), the following equation is obtained for the stress intensity factor at the beginning of service:

$$K_i < K_{TC} \sigma_a / \sigma_p$$
 (8)

Thus by proof testing a component prior to service, an estimate of the maximum stress intensity factor at the beginning of service may be obtained.

In service, the component will be gradually weakened by crack growth until failure occurs. The time to failure can be estimated from the definition of crack velocity, v=da/dt, and the relation between  $K_{\underline{I}}$ , a, and  $\sigma$ :  $(K_{\underline{I}}=\sigma Y\sqrt{a})$ . For a constant load

$$v = (2K_T/Y^2\sigma_a^2) dK_T/dt$$
 (9)

This equation can be solved to give the following expression for the failure time:

$$t = (2/Y^2 \sigma_a^2)$$
  $K_I^{C}$   $(K_I/v)$   $dK_I$  (10).

This equation can be evaluated either numerically or analytically if a relationship between  $K_{\underline{I}}$  and v is available. For  $v = AK_{\underline{I}}^n$ , a minimum time-to-failure,  $t_{\min}$ , is obtained by substituting  $K_{\underline{I}}$  from equation 8 into the lower integration limits of Eqn. (10).

$$t_{\min} = 2\sigma_a^{-2} (K_{IC} \sigma_a/\sigma_p)^{2-n}/AY^2(n-2)$$
 (11)

Note that all the variables in the right hand side of this equation can be determined experimentally.

The main features of Eqn. (11) is that the minimum time to failure  $t_{\min}$ , after a proof test is inversely proportional to the service stress squared and directly proportional to a function of the proof test ratio,  $\sigma_{\rm p}/\sigma_{\rm a}$ . In functional form the equation may be expressed as follows:

$$t_{\min} = \sigma_a^{-2} \cdot f(\sigma_p/\sigma_a)$$
 (12)

The function of proportionality,  $f(\sigma_p/\sigma_a)$ , is determined from measurements of  $K_{\rm IC}$  and from crack growth data. Thus, all the necessary information for failure prediction can be obtained by fracture mechanics techniques.

The relationship given by Eqn. (12) can be presented graphically in the form of a proof test diagram that can be used for design purposes. A diagram of this type is given in Fig. 8 for a low expansion silica glass

containing 7.5 percent TiO2. The diagram gives a plot of the minimum time to failure versus the service stress. Each of the lines on this diagram represents the relationship between  $t_{min}$  and  $\sigma_a$  for a given proof test ratio,  $\sigma_p/\sigma_a$ . The diagram is used to select an appropriate proof test ratio for a specific application. The silica glass represented by this diagram was considered for use as windows in the Space Shuttle, for which the glass was required to sustain a load of 3 MN/m<sup>2</sup> ( $\sim 4,000$  psi). Assuming that for safe operation the load would have to be supported for one year, a proof test ratio of approximately 2.6 is obtained from the diagram. Thus reliable performance

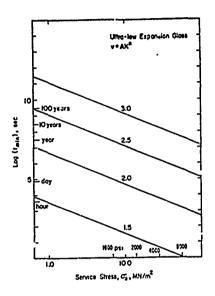


Fig. 8 Proof test diagram (Ref. 47)

of the space craft windows would require proof testing to a stress of 12  $\text{MN/m}^2$  (010,000 psi).

The above example of proof testing illustrates the basic procedure that should be followed to assure the reliability of glass for structural application. However, certain precautions are necessary to assure the validity of the method. These precautions are concerned with the accuracy and applicability of the proof test data, the design of a proper proof test procedure and the possibility of strength degradation after the proof test has been completed. These precautions are necessary since an incorrect assessment of any one of them could invalidate the proof test and result in premature component failure. The interested reader is referred to references 52-56 for a more complete discussion of the method and its application to structural design. Despite these precautions, proof testing is a promising method of using fracture mechanics data to assure the reliability of glass structures.

# SUMMARY

This paper has summarized the strength and fracture data on glass, emphasizing the phenomenon of static fatigue. In summarizing the strength of glass, it was noted that glass is characterized by extreme brittle behavior. The practical strength is much less than the theoretical strength because of the presence of cracks in the glass surface. The growth of these cracks caused by water in the environment leads to the phenomenon of static fatigue, which is characterized by delayed failure under constant load, or a loading rate dependence of strength.

Static fatigue can be studied by the techniques of fracture mechanics which use specimens that contain large artificially introduced cracks. The key variable that can be measured by fracture mechanics techniques is the stress intensity factor,  $K_{\rm I}$ , which is proportional to the stresses near the crack tip. By relating  $K_{\rm I}$  to the crack velocity, temperature, environment and glass composition it has been demonstrated that the fracture of glass is a complex phenomenon that is controlled by a number of different mechanisms involving chemical reactions, diffusion and glass structure. In cases where a direct comparison has been made between crack growth data and strength data, good agreement is obtained suggesting that crack growth data can be used for design purposes. A method of using fracture mechanics data to obtain proof test diagrams is outlined. These diagrams can be used to select the proof test load that will assure the lifetime of a structural component for a particular application.

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# REFERENCES

- (1) Hillig, W. B., C. R. Symposium sum la résistance mécanique du verre et les moyens de l'améliorer, Union Scientifique Continentale du Verre, Charleroi (1962).
- (2) Mallinder, F. P., Proctor, F. P., Phys. Chem. Glasses 5, 91 (1964).
- (3) Bridgman, P. W., and Simon, I., J. Appl. Phys. 24, 405 (1953).
- (4) Cohen, H. M., and Roy, R., Phys. Chem. Glasses 6, 149 (1965).
- (5) Ernsberger, F. M., J. Am. Ceram. Soc. 51, 545 (1968).
- (6) Dick, E. and Peter, K., J. Am. Ceram. Soc. 52, 338 (1969).
- (7) Ernsberger, F. M., Phys. Chem. Glasses 10, 240 (1969).
- (8) Ernsberger, F. M., Prog. Ceram. Sci. 3, 58 (1963).
- (9) Grenet, L., Bull. Soc. Encour. Ind. Mat. 4, 839 (1899).
- (10) Shand, E. B., J. Am. Ceram. Soc., 37, 52 (1954).
- (11) Milligan, L. H., J. Soc. Glass Technol., 13, 341T (1929).

- (12) Gurney, C. and Pearson, S., Proc. Phys. Soc., 62, 469 (1949).
- (13) Preston, F. W. and Baker, T. C., J. Appl. Phys. 17, 170 (1946).
- (14) Preston, F. W. and Baker, T. C., J. Appl. Phys. 17, 179 (1946).
- (15) Kropschot, R. H. and Mikesell, R. P., J. Appl. Phys. 28, 610 (1951)
- (16) Charles, R. J., Prog. Ceram. Sci. 1, 1 (1961).
- (17) Charles, R. J., J. Appl. Phys. 29, 1554 (1958).
- (18) Mould, R. E. and Southwick, R. D., J. Am. Ceram. Soc. 42, 489 (1959).
- (19) Holland, A. J. and Turner, W. E. S., J. Soc. Glass Technol. 24, 46T (1940).
- (20) Griffith, A. A., Phil, Trans. Ray. Soc. (London) 221A, 163 (1920).
- (21) Irwin, G. R., Encyclopedia of Physics, Springer, Heidelberg (1958).
- (22) Paris, P. C. and Sih, G. C., ASTM Special Technical Publication, No. 381.
- (23) Brown, W. F. and Srawley, J. E., ASTM Special Technical Publication No. 410
- (24) Srawley, J. E. and Gross, B., Mater, Res. Stand. 7, 155 (1967).
- (25) Wiederhorn, S. M., Shorb, A. M. and Moses, R. L., J. Appl. Phys. 39, 1569 (1968).
- (26) Evans, A. G., J. Mater. Sci. 7, 1137 (1972).
- (27) Evans, A. G., Fracture Mechanics of Ceramics, Plenum Press, New York (1974).
- (28) Wiederhorn, S. M., ibid.
- (29) Wiederhorn, S. M., J. Am. Ceram. Soc. 59, 407 (1967).
- (30) Wiederhorn, S. M. and Bolz, L. H., J. Am. Ceram. Soc. 53, 543 (1970).
- (31) Charles, R. J. and Hillig, W. B., C. R. Symposium sur la résistance mécanique du verre et les moyens de l'améliorer, Union Scientifique Continentale du Verre, Charleroi (1962).
- (32) Hillig, W. B. and Charles, R. J., High Strength Materials, John Wiley and Sons, New York (1965).
- (33) Wiederhorn, S. M and Johnson, H., J. Am. Ceram. Soc. <u>56</u>, 192 (1973).
- (34) Freiman, S. W., J. Am. Ceram. Soc., August (1974).
- (35) Evans, A. G. and Wiederhorn, S. M., J. Am. Ceram. Soc., To be published.
- (36) Wiederhorn, S. M., Johnson, H., Diness, A. M. and Heuer, A. H., J. Am. Ceram. Soc., July (1974).
- (37) Thomson, R., Hsleh, C. and Rana, R., J. Appl. Phys. 42, 3154 (1971).
- (38) Hsieh, C. and Thomson, R., J. Appl. Phys. 44, 2051 (1973).
- (39) Wiederhorn, S. M., J. Am. Ceram. Soc. <u>52</u>, 99 (1969).
- (40) Wiederhorn, S. M. Evans, A. G. and Roberts, D. E., Fracture Mechanics of Ceramics, Plenum Press, New York (1974).

- (41) Dugdale, D. S., J. Mech. Phys. Solids 8, 100 (1960).
- (42) McClintock, F. A. and Irwin, G. R., ASTM Special Technical Publication, No. 381.
- (43) Marsh, D. M., Proc. Ray. Soc. A279, 420 (1964).
- (44) McClintock, F. A. and Argon, A. S., Mechanical Behavior of Materials, Addison-Wesley Publication Co., Inc., Reading, Mass. (1966).
- (45) Evans, A. G., Int. J. Fract. Mech., To be published.
- (46) Evans, A. G. and Johnson, H., J. Mat. Sci,. To be published.
- (47) Wiederhorn, S. M., Evans, A. G., Fuller, W. R., and Johnson, H., J. Am. Ceram. Soc., August (1974).
- (48) Ritter, J. E. and Sherburne, C. L., J. Am. Ceram. Soc., 54, 601 (1971).
- (49) Ritter, J. E. and Manthuruthil, J., Glass Tech. 14, 60 (1973).
- (50) Doremus, R. H., Corrosion Fathque: Chemistry Mechanics and Microstructure, Nat. Ass. Corrosion Engr. (1972).
- (51) Wiederhorn, S. M and Johnson, H., unpublished data.
- (52) Evans, A. G. and Wiederhorn, S. M., Int. J. Fract. Mech. (1974), NBS Report No. NBSIR 73-147.
- (53) Wiederhorn, S. M., J. Am. Ceram. Soc. 56, 27 (1973).
- (54) Evans, A. G. and Fuller, E. R., Met. Trans. 5, 27 (1974).
- (55) Evans, A. G. and Fuller, E. R., J. Mech. Phys. Solids, To be published.
- (56) Wiederhorn, S. M., To be published in Proceedings of Ceramics for High Performance Application, Hyannis, Nov. 13-16 (1973).